

# Cyclotron resonance for 2D electrons on thin liquid helium films

A. Würfl<sup>†</sup>, V. Shikin<sup>‡</sup>, J. Klier<sup>†</sup> und P. Leiderer<sup>†</sup>

<sup>†</sup> Universität Konstanz, Germany; <sup>‡</sup> ISSP Chernogolovka, Moscow District, Russia

## Introduction & Motivation

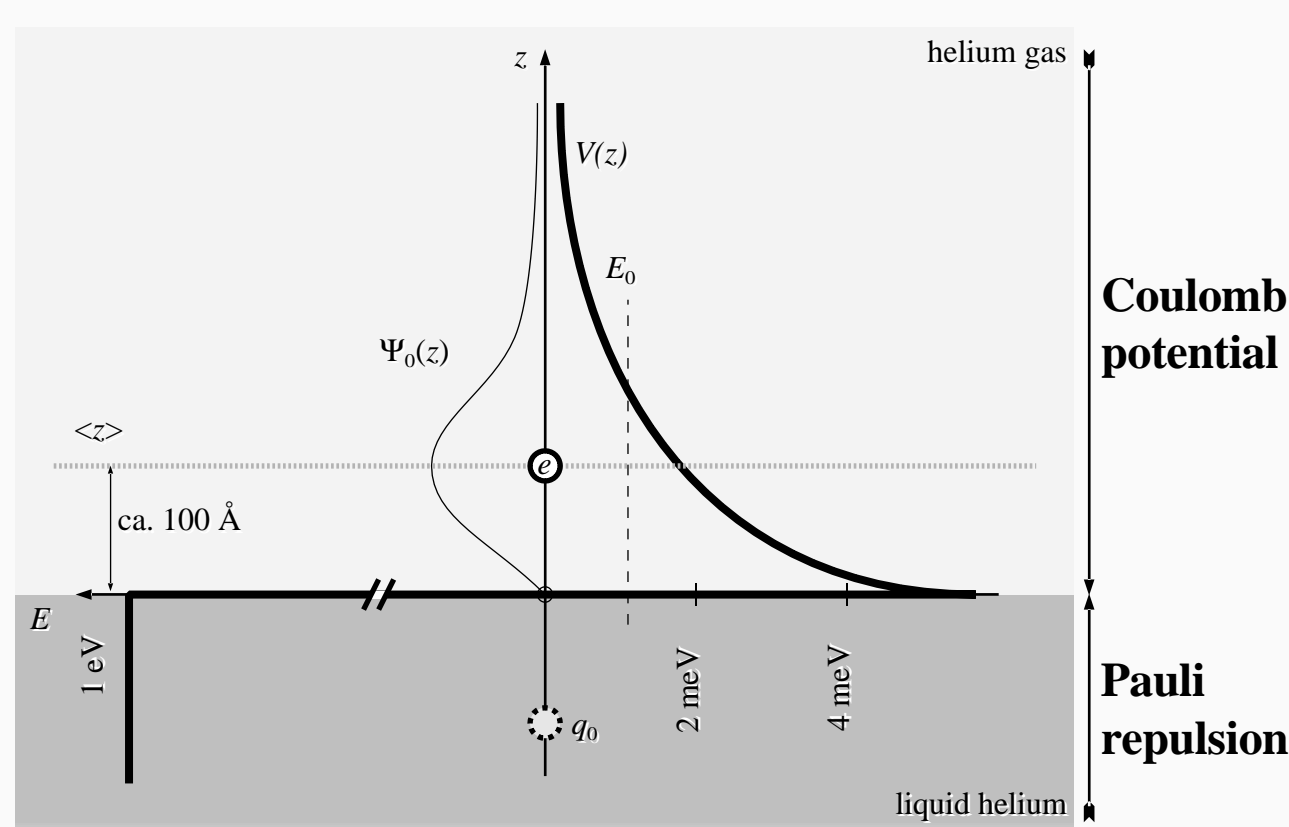
Cyclotron resonance (CR) measurements of two-dimensional electron systems (2DES) on thin helium films show an increasing asymmetry of the CR-line when the helium film thickness is reduced. So far there was no satisfying explanation of this effect. We are presenting a new theoretical approach taking the roughness of the substrate-surface into account. On a rough surface **two fractions** of electrons occur, free electrons and electrons localized to tops of the surface roughness. The well known CR-line for free electrons is then modified by the contribution of the localized electrons, which gives rise to the asymmetry.

## Electrons on liquid Helium

Electrons on liquid helium feel a hydrogen-like potential

$$V(z) = \begin{cases} V_0 & z \leq 0 \\ -\frac{1}{4\pi\epsilon_0} \frac{q_0 e^2}{z+\beta} & z > 0 \end{cases} \quad \text{with:} \quad \begin{cases} V_0 \approx 1 \text{ eV} \\ q_0 = \frac{\epsilon_L(\text{He})-1}{4(\epsilon_L(\text{He})+1)} \end{cases}$$

If one assumes  $V_0$  to be infinite, the resulting wavefunctions are similar

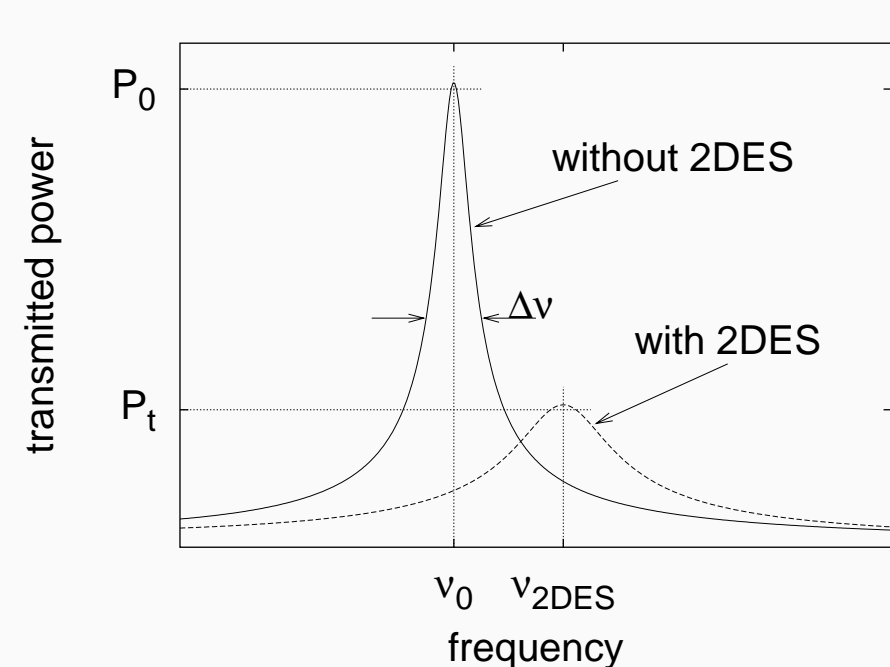
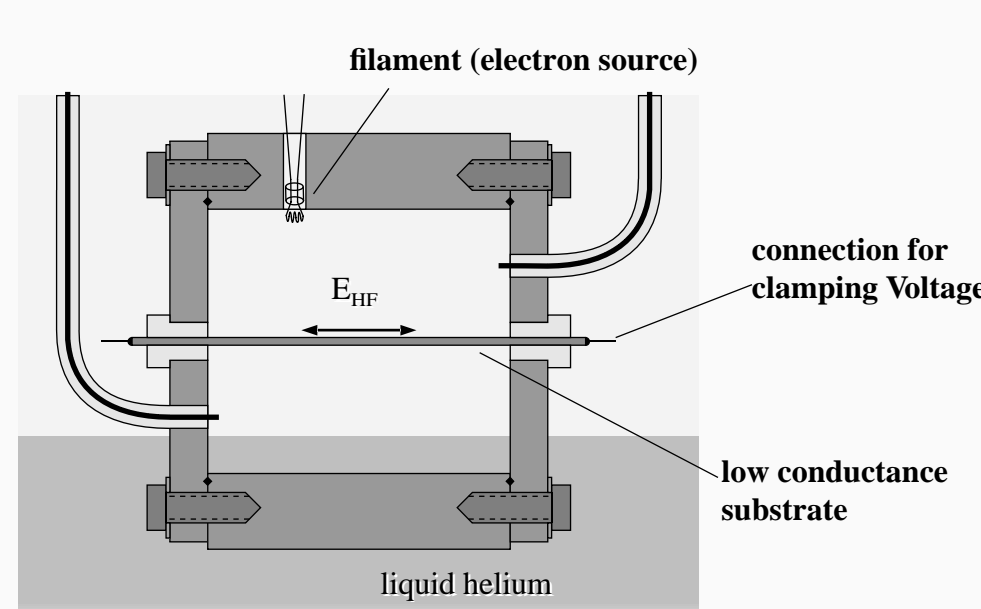


to the well-known wavefunction of the hydrogen atom.

In the experiment,  $T$  is around 1.3K. At this temperature most of the electrons are in the ground state and therefore form a real 2D layer on top of the liquid helium.

## Experimental Setup

A schematic drawing of the used microwave cavity is shown. It is driven in the fundamental TM010 mode, where the maximal amplitude of the electrical field is in the vicinity and parallel to the substrate.



The dependence of the resonance of the cavity with and without a 2DES on the helium film on the substrate is shown.

## The two-fraction Model

### Modelling the Helium film thickness

For structures much smaller than the capillary length of liquid helium one can extract the properties of a neutral helium film  $d(x)$  from

$$\sigma_{lv} \frac{d''(x)}{[1 + (d')^2]^{3/2}} - \rho g \delta(x) + \frac{C_3}{d^3(x)} = \rho g h$$

where  $C_3$  is the van-der-Waals constant of the helium-substrate boundary. The radius of curvature of the capillary condensed film is defined as

$$\frac{2\sigma_{lv}}{R} \approx \rho g h$$

### Modelling the substrate roughness

We assume that a 1-dimensional roughness behaviour  $\delta(x)$  can be described by a Gaussian distribution of the amplitudes

$$G(\delta) = \frac{1}{\sqrt{2\pi}\Delta^2} \exp\left[-\frac{\delta^2}{2\Delta^2}\right]$$

and by the lateral correlation length

$$\langle \delta(x), \delta(x-x') \rangle = \Delta^2 \exp\left(-\frac{x'^2}{2\eta^2}\right).$$

Here only the high enough tops of the substrate roughness above a the fixed level  $\delta > 0$  play a role. Their density is given by

$$n_\delta = \frac{1}{\sqrt{2\pi}\eta} \exp\left(-\frac{\delta^2}{2\Delta^2}\right) \quad \text{with} \quad \eta = \sqrt{\langle \eta^2 \rangle}$$

This gives, with some limitation, the density of the active tops  $n_a^T$ .

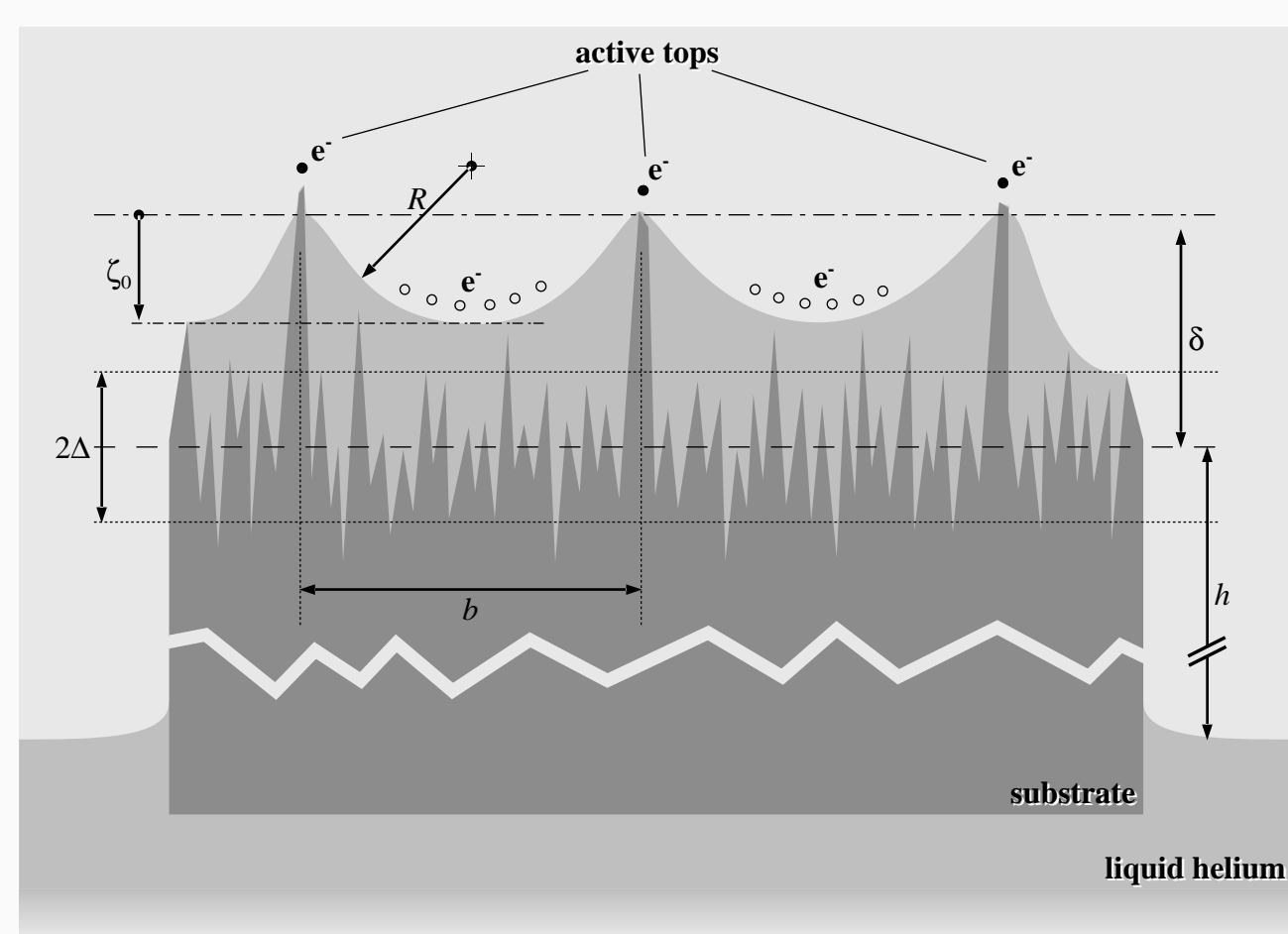
### Properties of the two fractions

The total electron density  $n_s$  is the sum of free and localized density

$$n_e + n_l = n_s$$

The relationship between these fractions is flexible and is determined by the chemical potential  $\mu_0$ . The definition of  $\mu_0$  is taken from semiconductor physics

$$n_l = \frac{n_a}{\exp((V_a - \mu_0)/T) + 1}, \quad V_a < 0, \quad n_e = \frac{n_a^e \exp(T_e/T)}{\exp(-\mu_0/T) + 1}, \quad n_e^e = \frac{mT}{(2\pi\hbar^2)^{3/2}}$$



### Results for CR in the microwave cavity

The absorption of the free electron fraction is

$$Q_e^{-1} \propto n_e \frac{1+z+x}{(1-z+x)^2 + 4z}$$

The absorption of the localized electron fraction  $n_l$  is

$$Q_l^{-1} \propto n_l \frac{\arctan \frac{\sqrt{z}}{1+x+\sqrt{xz}} + \arctan \frac{\sqrt{z}}{(1+x)\sqrt{z}-z\sqrt{x}} + c(z, x)}{2\sqrt{z}}$$

Here  $z = \omega_0^2 \tau^2$  and  $x = \omega_c^2 \tau^2$ . So the total absorption of the 2DES in an external magnetic field is given by the sum of its parts:

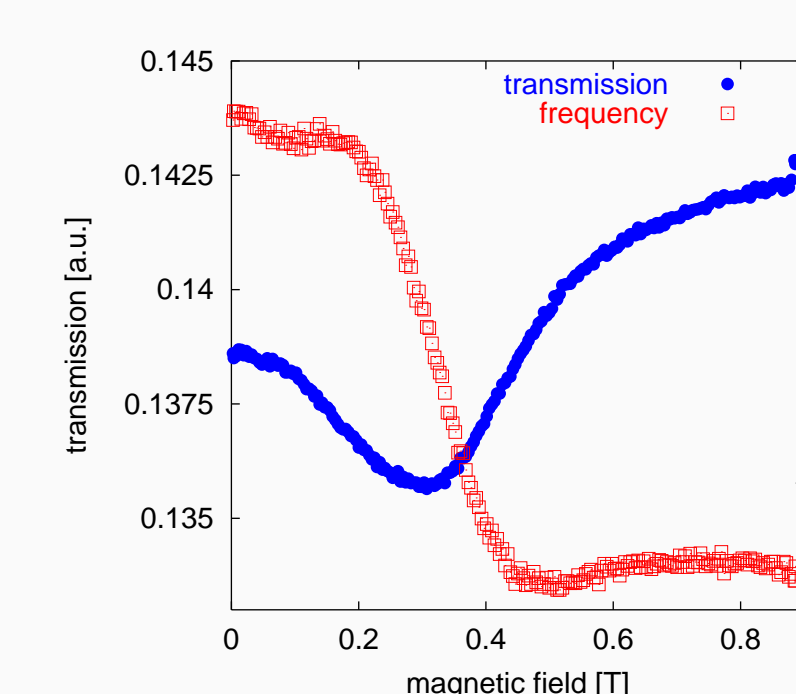
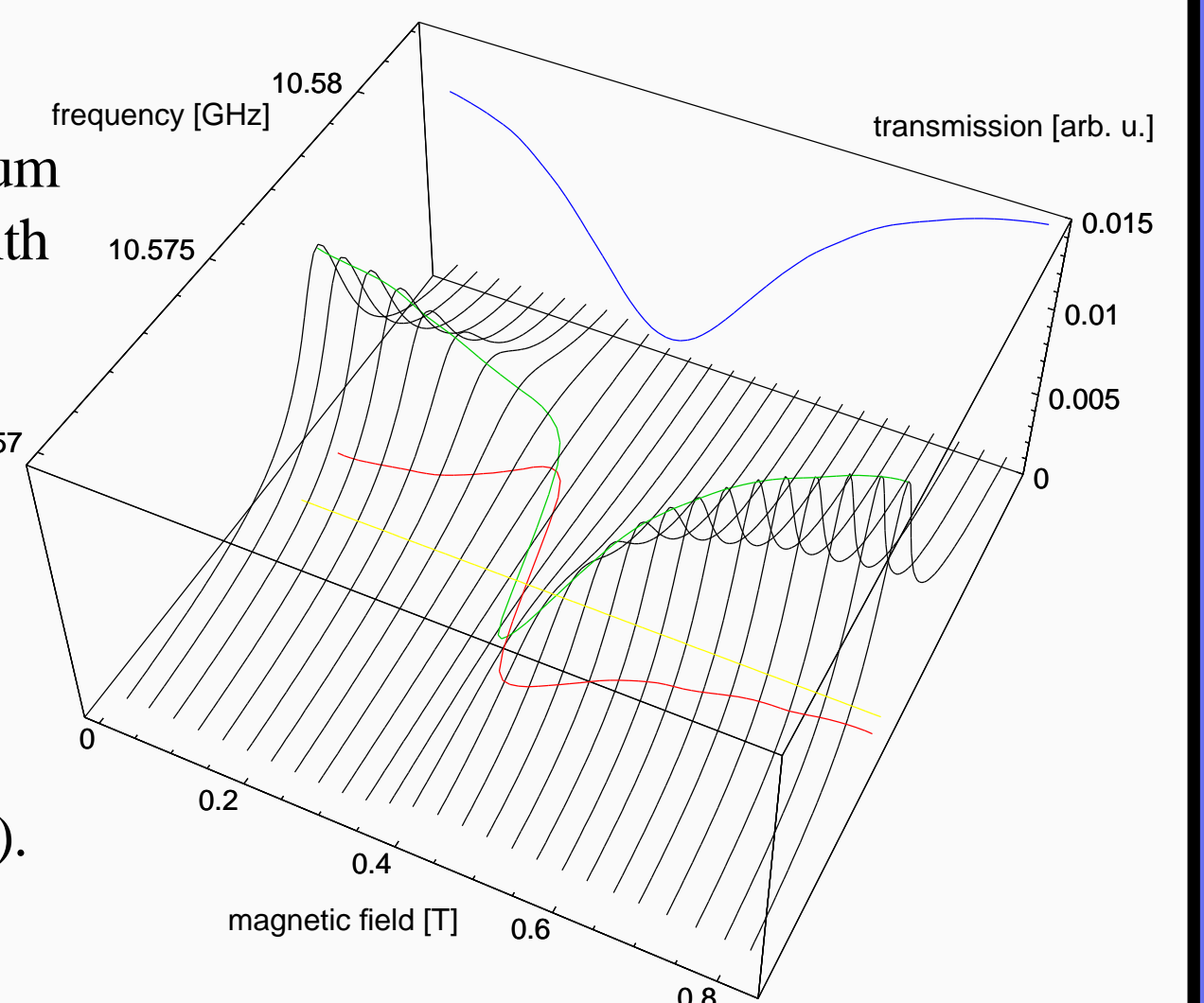
$$Q^{-1} = Q_e^{-1} + Q_l^{-1}$$

## Experiments

### How the experiment works

After charging the helium film on the substrate with a 2DES, the magnetic field is swept from 0 to around 0.9 T.

The resonance of the cavity is recorded for different values of the magnetic field. (Here for a 2DES on bulk He).



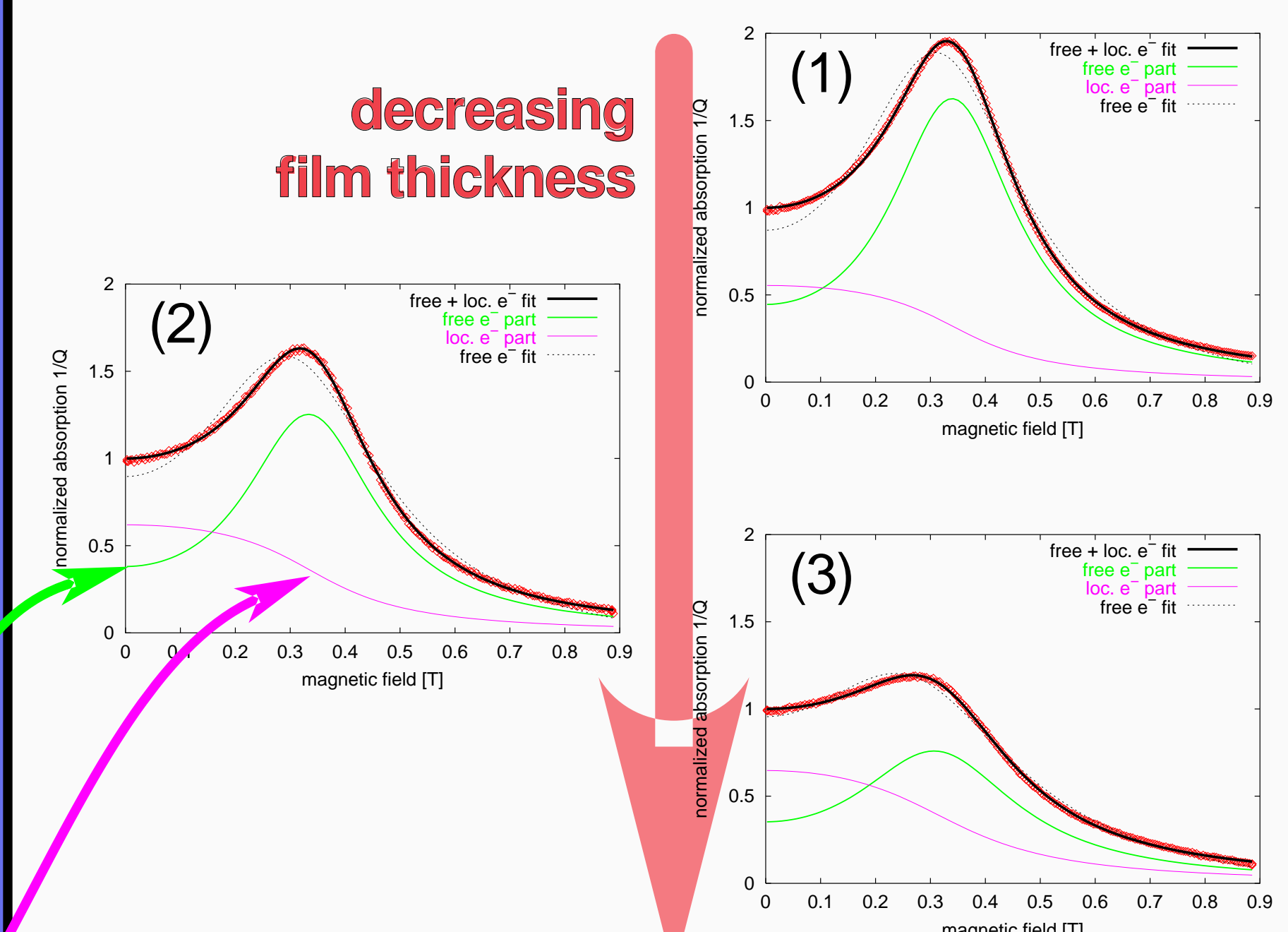
The position of the resonance maximum in transmission and frequency gives the typical CR data.

### Fitting the results

The reciprocal of the experimental transmission data is fitted to  $Q^{-1}$ , with the following 5 free parameters:

$\omega$ ,  $\tau$ ,  $A_{\text{free}}$  electrons,  $A_{\text{localized}}$  electrons and  $c$  constant offset

Here the fit results, normalized to  $Q^{-1}(B=0 \text{ T})$ , are shown. The green and the pink lines represent the **free** and **localized** electron fraction, the **full** line is the sum of both, fitted to the data:



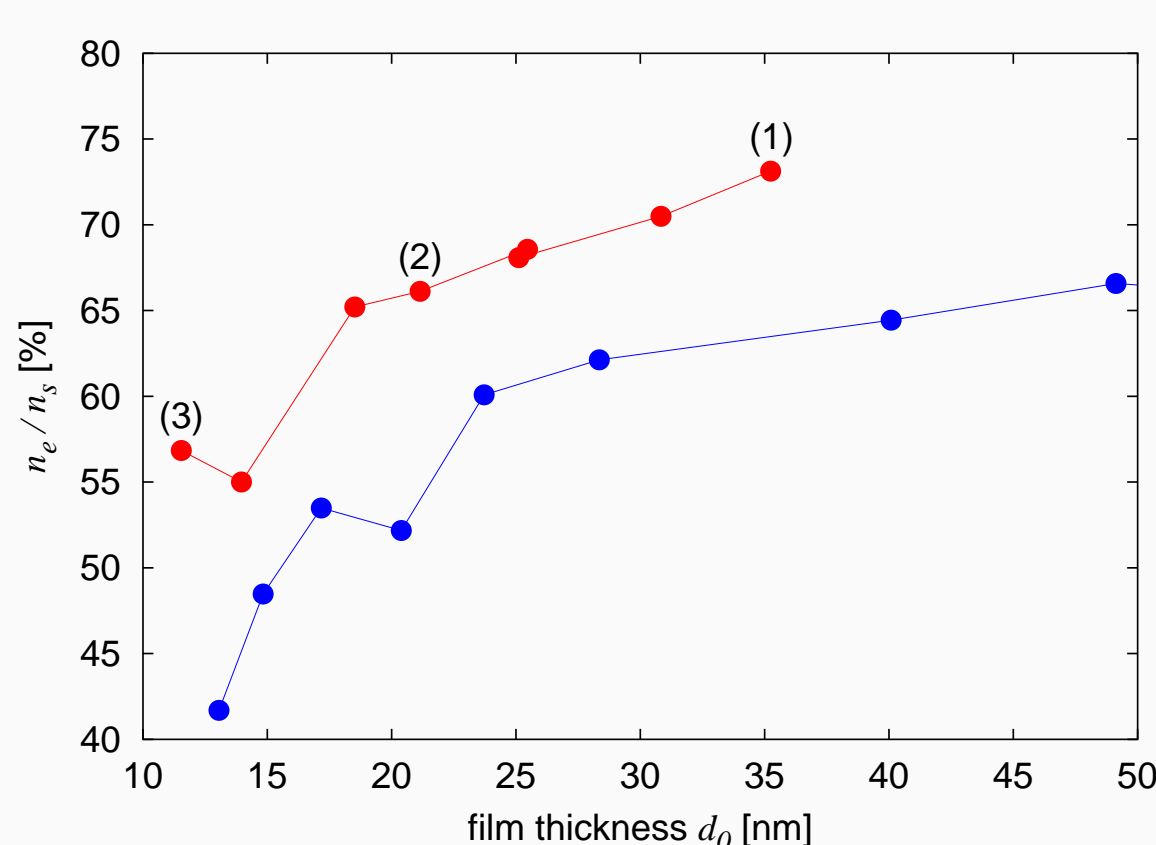
To obtain the values of the

$$\text{free } \nu_e = \frac{n_e}{n_s} \text{ and localized } \nu_l = \frac{n_l}{n_s} \text{ electron fraction,}$$

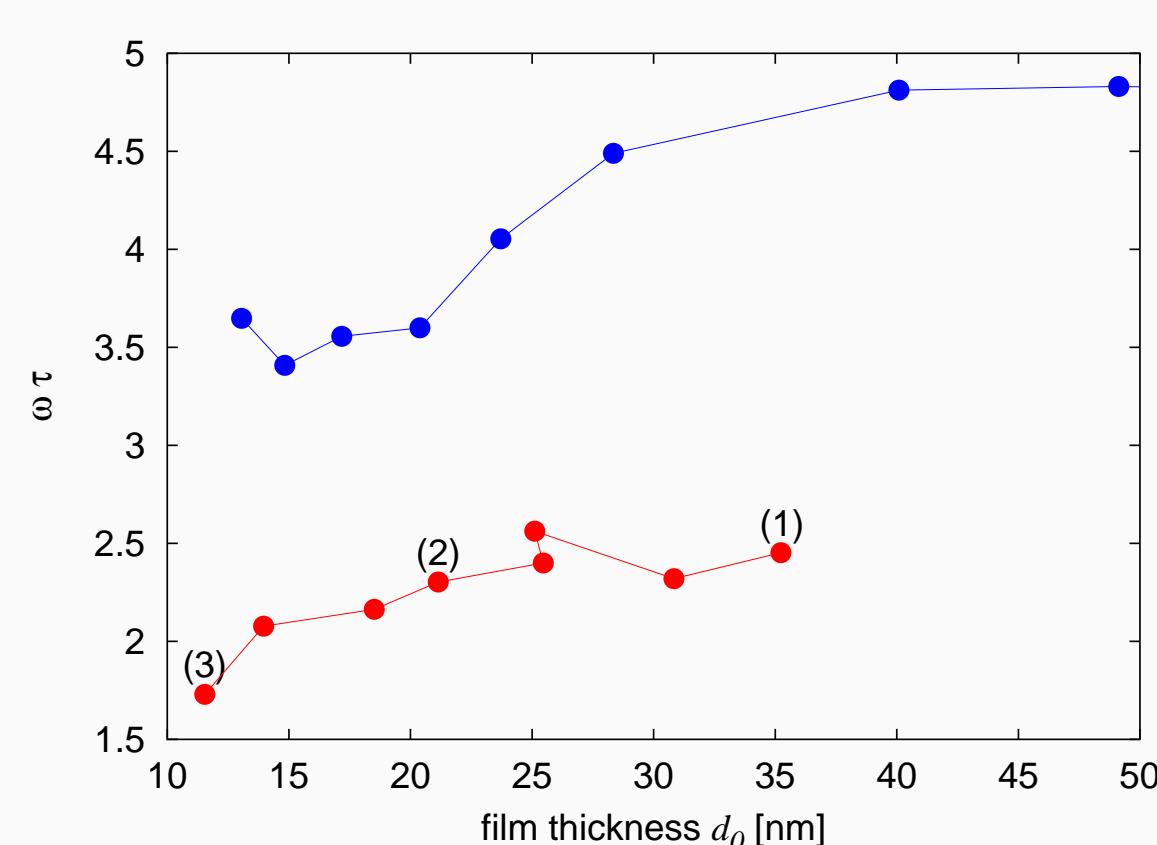
the ratio between the maximum and the zero field value of  $Q^{-1}$  is used:

$$\frac{Q^{-1}(\omega_c^{(max)})}{Q^{-1}(\omega_c = 0)} = \frac{\nu_e p(\omega_0, \tau, \omega_c^{(max)}) + \nu_l q(\omega_0, \tau, \omega_c^{(max)})}{\nu_e p(\omega_0, \tau, 0) + \nu_l q(\omega_0, \tau, 0)}$$

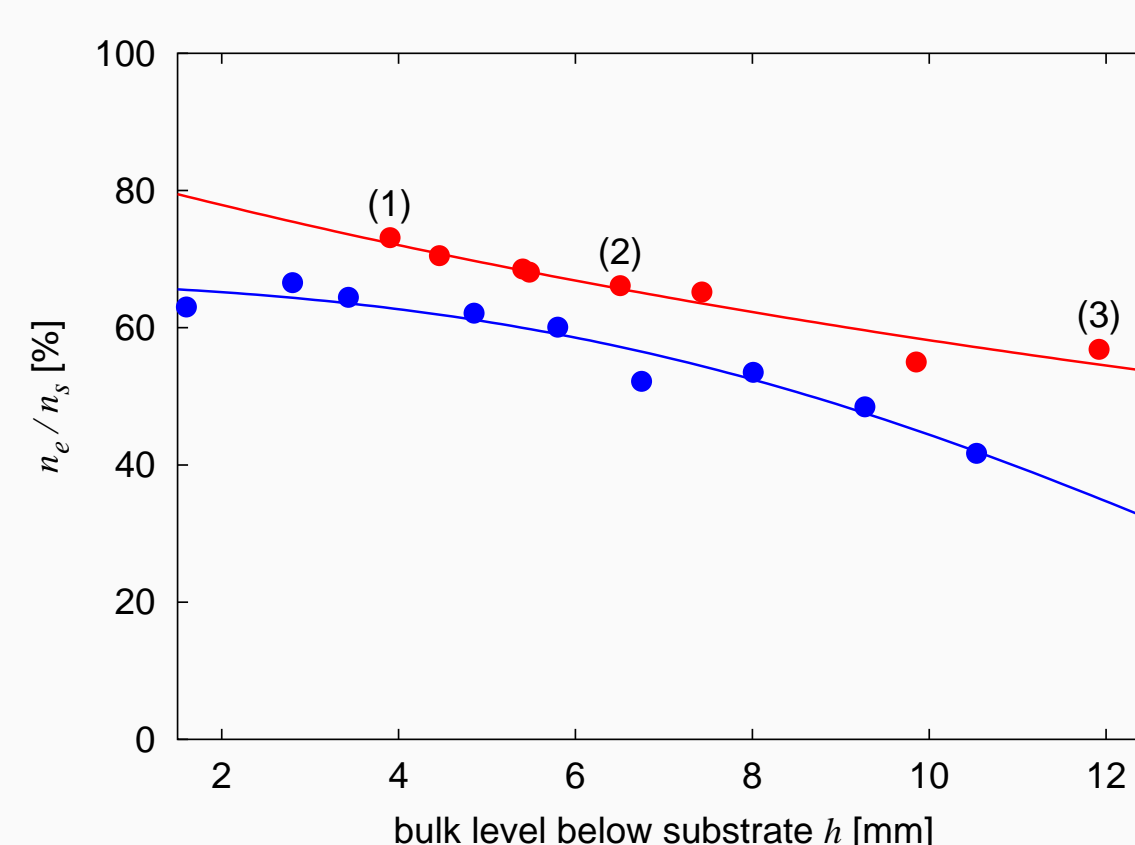
## Analyzing the experimental data



The free electron fraction vs thickness of the helium film shows qualitatively the expected behaviour. The red data-set is measured on a very smooth **SiO<sub>2</sub> on Silicon** substrate. The blue data-set is taken on **PMMA on Silicon** substrate (by G. Mistura). The dip in the thin film region may be related to the so-called dip-effect, described by *Shikin et al., PRB 64 (2001)*



The behaviour of  $\omega\tau$  vs the thickness of the helium film is, at least qualitatively as expected. Unfortunately the presented theory is too crude to provide a detailed interpretation.



Fit of the model parameters of the localized electrons to the data. From the data on PMMA we get  $\Delta$  around 8 nm and  $\eta$  around 6 nm. These values, being in the nm-regime, are reasonable and typical for the surfaces used in the experiments.

## Conclusions

The two-fraction model explains the behaviour of electrons on He films and is the basis for:

- the origin of CR-line asymmetry of a 2DES on thin helium films
- a quantitative analysis of the electron fraction
- the characterization of the substrate surface

## Outlook

Experimentally:

- Perform CR measurements for a wider range of film thickness and substrate roughness.
- Do precise measurements to resolve the dip and its origin in the  $n_e/n_s$  vs  $d_0$  diagram

Theoretically:

- Refine theory for quantitative data analysis.
- Consider special properties of the cavity.

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